



This is a digital copy of a book that was preserved for generations on library shelves before it was carefully scanned by Google as part of a project to make the world's books discoverable online.

It has survived long enough for the copyright to expire and the book to enter the public domain. A public domain book is one that was never subject to copyright or whose legal copyright term has expired. Whether a book is in the public domain may vary country to country. Public domain books are our gateways to the past, representing a wealth of history, culture and knowledge that's often difficult to discover.

Marks, notations and other marginalia present in the original volume will appear in this file - a reminder of this book's long journey from the publisher to a library and finally to you.

### Usage guidelines

Google is proud to partner with libraries to digitize public domain materials and make them widely accessible. Public domain books belong to the public and we are merely their custodians. Nevertheless, this work is expensive, so in order to keep providing this resource, we have taken steps to prevent abuse by commercial parties, including placing technical restrictions on automated querying.

We also ask that you:

- + *Make non-commercial use of the files* We designed Google Book Search for use by individuals, and we request that you use these files for personal, non-commercial purposes.
- + *Refrain from automated querying* Do not send automated queries of any sort to Google's system: If you are conducting research on machine translation, optical character recognition or other areas where access to a large amount of text is helpful, please contact us. We encourage the use of public domain materials for these purposes and may be able to help.
- + *Maintain attribution* The Google "watermark" you see on each file is essential for informing people about this project and helping them find additional materials through Google Book Search. Please do not remove it.
- + *Keep it legal* Whatever your use, remember that you are responsible for ensuring that what you are doing is legal. Do not assume that just because we believe a book is in the public domain for users in the United States, that the work is also in the public domain for users in other countries. Whether a book is still in copyright varies from country to country, and we can't offer guidance on whether any specific use of any specific book is allowed. Please do not assume that a book's appearance in Google Book Search means it can be used in any manner anywhere in the world. Copyright infringement liability can be quite severe.

### About Google Book Search

Google's mission is to organize the world's information and to make it universally accessible and useful. Google Book Search helps readers discover the world's books while helping authors and publishers reach new audiences. You can search through the full text of this book on the web at <http://books.google.com/>

QE  
433  
D4

UC-NRLF



\$B 51 030

YC 40243

Digitized by Google

LIBRARY  
OF THE  
UNIVERSITY OF CALIFORNIA.

GIFT OF  
*Univ. of Pennsylvania*  
Class













# A METHOD OF PETROGRAPHIC ANALYSIS

BASED UPON

## Chromatic Interference with Thin Sections of Doubly-Refracting Crystals in Parallel Polarized Light.

### Thesis

Presented to the Faculty of Philosophy of the University of Pennsylvania in partial fulfillment of the requirements for the degree of Doctor of Philosophy.

BY

HOMER MUNRO DERR, A. M.

Hector Tyndale Fellow.

The Randal Morgan Laboratory of Physics.

1903.

PLATE 1.

**Thin Section showing Free Gold in Metamorphic Diorite  
FROM ENCAMPMENT, WYOMING.**

Magnified 50 diameters:—*f*, feldspar, (oligoclase, albite); *q*, quartz;  
*h*, hornblende; *m*, magnetite; *g*, gold.  
Crossed nicols; thickness of section, 56 microns.







## BEHAVIOR OF ISOMETRIC CRYSTALS.

Isometric crystals on account of their isotropic character exhibit no special phenomena in polarized light. If a thin plate of such a mineral be placed in the path of a polarized beam of light between the polarizer and the analyzer, the beam will experience no alteration of its plane of vibration, no matter in what direction the plate was cut from the crystal, nor in what position it lies between the polarizer and the analyzer ; for the elasticity of the ether is the same in all directions in such a case, and rotating the mineral about any axis whatever will effect no change of the plane of vibration of the polarized light. If the section is colorless, it will not alter the color or the brightness of the field of view in a polarization microscope, except for the small amount of absorption which a ray of light will experience in passing through any medium. If the section is colored, the field of view will appear slightly different in color from that of the mineral ; however, this color does not change by altering the position of the plate in any way. Furthermore, when the nicols are crossed, a section on the stage of the microscope appears dark, and revolving it in any plane produces no effect. It appears light, and experiences no change in any position between parallel nicols. Sections of transparent isometric crystals may always be recognized as such by the fact that they behave as an amorphous substance in polarized light. There are, however, some optical anomalies in this connection. Sections of amorphous or isometric minerals sometimes appear light when viewed between crossed nicols. These results are due to internal strains caused either by inclusions of gases or fluids which exert a pressure on their surroundings, or by contraction of adjacent parts. These phenomena may usually be distinguished from ordinary double refraction, since the appearance is not generally uniform throughout the section.

## INTERFERENCE.

Fresnel and Arago investigated completely the conditions of interference of two rays polarized at right angles to each other

after they had been brought back to the same plane of polarization by passing them through a crystal of calc-spar whose principal section made an angle of  $45^\circ$  with the planes of polarization of each of the two rays. They found, (1) that two rays polarized at right angles to each other, which have come from an unpolarized ray, do not interfere even when they are brought into the same plane of polarization; (2) that two rays polarized at right angles, which have come from a polarized ray, interfere when they are brought back to the same plane of polarization.

Let a ray of polarized monochromatic light be incident at the lower surface of a plane-parallel doubly-refracting plate at any angle, Fig. 1, at the point A. It is separated into two rays AD and AE, vibrating in planes at right angles to each other and following different paths in the plate. On emergence they follow paths parallel to their direction at incidence, but are not coincident, and do not produce interference. But other parallel rays from the same source and incident at B and C, will emerge from

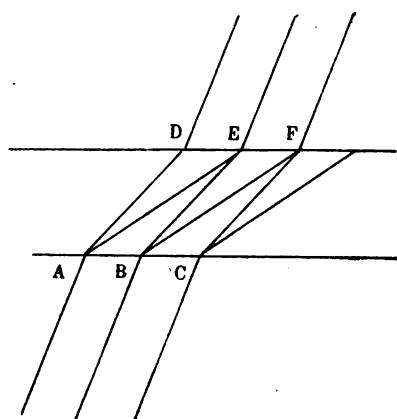


FIG. 1.

the upper surface, so that from all points E and F, the ordinary component of one ray and the extraordinary of another follow the same path. These rays will have travelled over slightly different paths in the plate, and with different velocities; then on emergence one ray must have advanced a certain number of wave-lengths, or fractions thereof, ahead of the other. The waves are therefore in different phases and retain this difference of phase while they travel in the same medium. The plane of vibration of the extraordinary ray is parallel to the principal optic section of the plate, and that of the ordinary ray is at right angles to the same. If an analyzer be placed in the path of these rays, each ray will be by it resolved into components the vibrations of which are in and at right angles to the principal section of the polarizer, and only

the former will be transmitted. That is to say, there will emerge two rays advancing in the same line and with parallel vibrations; hence they are capable of interfering, since light-rays can completely interfere only when their vibration are in the same plane. If these vibrations are in the same phase, the intensity of the resultant wave will be proportional to the square of the sum of their amplitudes; but if in opposite phases, the intensity will be proportional to the square of the difference.

## WITH PARALLEL MONOCHROMATIC LIGHT AND CROSSED NICOLS.

In sections cut perpendicular to the optic axis, interference phenomena are impossible, and the field remains dark throughout a rotation of  $360^\circ$ , since the light from the polarizer traverses the section in the direction of the optic axis, and therefore without change. In all other sections there is double refraction, and consequently interference. The field is dark four times during the entire rotation of the stage at intervals of  $90^\circ$ ; that is, when the planes of vibration of the rays produced in the section coincide with the principal sections of either nicol. For all other positions the field is illuminated by the components of the rays which penetrate the analyzer, and this brightening is most intense in the diagonal positions.

The rays pursuing the same path are brought into one plane of vibration by the analyzer and there interfere. The kind of interference is determined by the formula,

$$\Delta = h \left[ n_1 - n_2 + \sin^2 i \left( \frac{1}{n_2} - \frac{1}{n_1} \right) \right],$$

in which  $\Delta$  is the difference in the retardations in microns which the two rays have undergone.

$h$  is the thickness of the section in microns.

$\lambda$  is the wave-length in microns.

$i$  is the angle of incidence.

$n_1$  is the index of refraction for the slower ray.

$n_2$  is the index of refraction for the faster ray.

For normal incidence, which is usually the case in practice,  
 $\Delta = h (n_1 - n_2)$ .

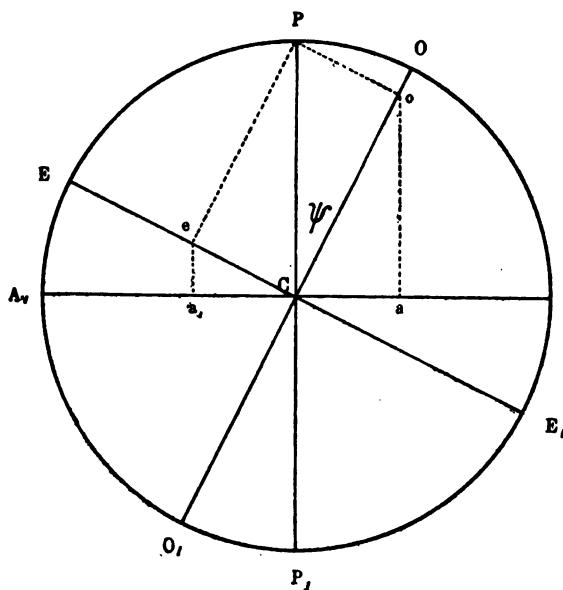


FIG. 2.

acting upon any ether particle C are  $C_O$  and  $C_E$ , which when reduced to the plane of the analyzer are  $C_a$  and  $C_{a_1}$ , equal and in opposite directions.

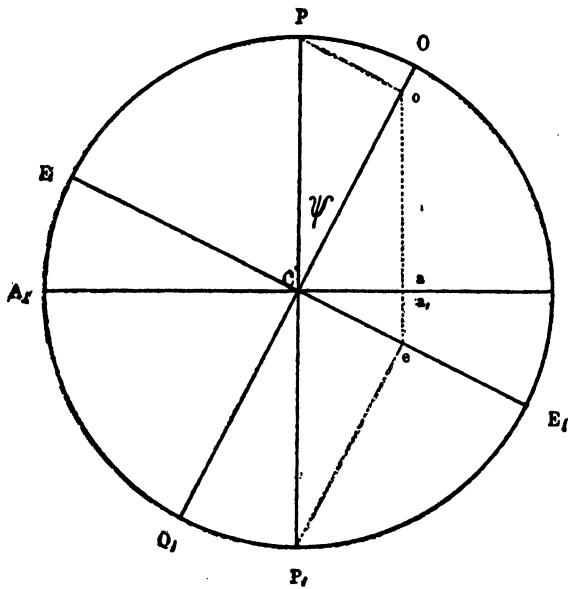


FIG. 3.

When  $\Delta = \lambda$ ,  $2\lambda$ ,  $3\lambda$ , etc., the field is dark during an entire revolution of the section on the stage. From Fig. 2, it will be seen that the components of  $PP_1$ , on emergence from the plate with vibration directions  $OO_1$  and  $EE_1$ , must be of the same phase; that is, the simultaneously displacing forces

When  $\Delta = \frac{1}{2}\lambda$ ,  $\frac{3}{2}\lambda$ ,  $\frac{5}{2}\lambda$ , etc., the light will be at its brightest because the components of  $PP_1$  must then be of opposite phase on emergence from the plate. From Fig. 3, the simultaneously displacing forces acting on any ether particle C are  $C_O$  and  $C_E$ , which reduced by the

analyzer to its plane are  $C_a$  and  $C_{a_1}$ , equal and in the same direction.

Let a thin wedge of doubly-refracting crystal, such as quartz, be cut so that its planes of vibration are parallel to the length and breadth, and then examined between crossed nicols, using perpendicularly incident monochromatic light. Upon rotation, it will appear dark when in the "normal positions," but in all others it will show a series of dark and light bands which are most marked in the diagonal position. When the nicols are made parallel the portions formerly light become dark, and vice versa. The distance between the dark bands varies with light of different wave-lengths.

The formula for intensity of the emerging light is

$$I = a^2 \sin^2 2 \Psi \sin^2 \left[ \frac{\pi \Delta}{\lambda} \right], *$$

in which  $a$  = the amplitude of the incident wave,  $\lambda$  = the wave-length,  $\Delta$  = the retardation,  $\Psi$  = the angle between the vibration plane of the polarizer and the slower wave.

I will be a minimum :

(a) when  $\sin^2 2 \Psi = 0$ , or  $2 \Psi = 0^\circ, 180^\circ, 360^\circ, 540^\circ$ , etc., or  $\Psi = 0^\circ, 90^\circ, 180^\circ, 270^\circ$ , etc. ; that is, four times in a revolution of the section, or whenever the planes of vibration of the plate coincide with those of the polarizer and the analyzer.

(b) when  $\sin^2 \left[ \frac{\pi \Delta}{\lambda} \right] = 0$ , which will be whenever  $\frac{\Delta}{\lambda} = 1, 2, 3$ , etc. ; that is, whenever the phase difference  $\Delta$  is a multiple of  $\lambda$ , for then  $\sin^2 \left[ \frac{\pi \Delta}{\lambda} \right]$  becomes  $\sin^2 180^\circ$ , or a multiple thereof. This is independent of  $\Psi$  ; hence with this condition, the section will appear dark throughout an entire revolution.

I will be a maximum :

When  $\sin^2 2 \Psi = 1$  or  $2 \Psi = 90^\circ, 270^\circ$ , etc., and  $\Psi = 45^\circ, 135^\circ, 225^\circ$ , etc.

When  $\frac{\Delta}{\lambda} = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}$ , etc., for then  $\sin^2 \left[ \frac{\pi \Delta}{\lambda} \right] = \sin^2 90^\circ, \sin^2 270^\circ$ , etc., = 1.

---

\* See Derivation of Formulae, page 21.

## WITH PARALLEL WHITE LIGHT AND CROSSED NICOLS.

In sections cut normal to the optic axis, the field appears dark throughout a rotation of  $360^\circ$ , the optic axis being the same for all colors. In all other sections there is extinction every  $90^\circ$ , and greatest brightness in the diagonal positions; but, since  $\Delta$  may be at the same time approximately an even multiple of  $\frac{1}{2}\lambda$  and an odd multiple of  $\frac{1}{2}\lambda^1$ , light of one wavelength may be greatly weakened while that of another wavelength is practically undimmed; that is, there will result a tint due to the partial extinction of some colors.

For example, suppose that in a polarization microscope, parallel white light passes upward through the polarizer, whose principal section is represented by  $PP_1$  in Fig. 2. Next let this light, which is polarized in a single plane  $PP_1$ , pass through a thin section of gypsum. It will there be separated into two waves vibrating in planes at right angles to each other,  $OO_1$  and  $EE_1$ . The two waves travel through the section with unequal velocity, and on emerging one is retarded a certain number of wave-lengths, or fractions thereof, as compared with the other. Now let these light waves pass through the analyzer, with its vibration-plane  $AA_1$  at right angles to  $PP_1$ . Then each of the two sets of vibrations will have a component in the direction  $AA_1$ , and these will emerge polarized in the same plane, and are therefore capable of interfering. Light corresponding to the other components in the direction  $PP_1$  will be extinguished. As previously stated, one of the emergent waves is slightly retarded as compared with the other. The amount of this retardation, on which the interference-color of the section depends, is proportional to the difference of the indices, and also to the thickness of the section.

If in the examination of the plate of gypsum above, one of the nicols had been rotated  $90^\circ$ ,—that is to say, if the principal sections of the two nicols had been parallel,—interference would have taken place between the emerging rays; but the color resulting in each case would have been exactly *complementary* to that obtained at first when the nicols were crossed. For instance, in Fig 2, if  $\Psi = 0$ , the principal sections of the plate coincide with those of polarizer and analyzer, and complete extinction occurs; but if  $AA_1$  be made parallel to  $PP_1$ , evidently the section would appear white.

If we examine in the same way a very thin and gradually-tapering wedge of some doubly-refracting crystal (quartz for example), cut so that its planes of vibration are parallel to the length and breadth, the successive interference-colors of the *first order*, beginning at the thin end, pass from an iron-gray through bluish-gray to white, yellow, and red; then follow violet, indigo, blue, green, yellow, orange, and red of the *second order*; then the similar but paler series of colors of the *third order*, and finally the very pale shades of green and red of the *fourth order*. Beyond this the colors are not very distinct, and white of a higher order finally results from the interference. A mineral of very strong double refraction, such as calcite, shows only the white of the higher orders unless extremely thin.

In general, a thin section of a doubly-refracting crystal, examined between crossed nicols, is not dark except when placed in certain definite positions. In any other position it does not completely extinguish the light, but its effect, in conjunction with the nicols, is partially to suppress the several components of the white light in different degrees, so that in the emergent beam these components are no longer in the proportions to give white light. In this way arise interference-tints, which may be definitely classed according to Newton's color-scale. (See page 15.) The several tints, though graduating into one another, are distinguished by names and divided into several orders. The precise position in the scale of a given tint observed between crossed nicols can be fixed by means of a quartz-wedge or other contrivance for compensating or neutralizing the effect of double refraction of the section.

The interference-tints given by a crystal section depend (1) on the amount of double refraction of the mineral ( $n_1 - n_2$ ), which is a specific character; (2) on the direction of the section relatively to the ellipsoid of optic elasticity, the tint being highest for a section parallel to the greatest and least axes of the ellipsoid; (3) on the thickness of the section. The last two are disturbing factors, which must be eliminated before we can use the interference-tints as an index of the amount of double refraction of the crystal, and as a useful criterion in identifying the mineral. The fact that the interference-tints depend in part on the direction of the section through the crystal will give little difficulty in estimating approximately the amount of double refraction of the mineral. If several crystals of the

same mineral are contained in a rock-section, it is sufficient to have regard to the ones which give the highest interference-tints. Even a single crystal will in the majority of cases give tints not so far below those proper to the mineral as to occasion error, but the possibility of the section having an unfavorable direction must be borne in mind.

In making a determination, the manner of procedure is briefly as follows: Select a crystal, revolve it on the stage of the microscope until it is in the position of maximum intensity (i. e., when the greatest and least axes of elasticity make an angle of about  $45^{\circ}$  with the principal sections of polarizer and analyzer); introduce the quartz-wedge between polarizer and analyzer (in slot provided for the purpose, also making an angle of  $45^{\circ}$  with principal sections of the two nicols), until a position of maximum darkness of the crystal is reached. If the interference-tint of the crystal cannot in this position, or when revolved through an angle of  $90^{\circ}$ , be reduced through the successive colors of the scale, in descending order, back to darkness, the section is an unfavorable one, i. e., it is not cut approximately parallel to the plane of the axes of greatest and least elasticity of the crystal. When a favorable section is found, after having fixed the quartz-wedge in the position giving maximum darkness to the crystal, first remove the section from the stage of the microscope. Now the quartz-wedge, as viewed through the microscope, will be of the same color and order in Newton's scale as that of the crystal-section under the same conditions. Carefully withdraw the wedge, at the same time noting the succession of tints in descending order. In this fashion, one can estimate the order of the interference-color in the scale; and by applying the formula  $\Delta = h(n_1 - n_2)$ , in connection with the tables given on pages 15 and 16, many of the commoner rock-forming minerals may be identified.

The following twenty-four examples will illustrate:

Rock Section.	Thickness in Microns.	Crystals.	Interference Color between Crossed Nicols.	Interference Color between Parallel Nicols.	Order of Color.
QUARTZ-PORPHYRY ...	26.5	Quartz Orthoclase	Greenish white Gray	Brown Brownish yellow	I I
RHYOLITE .....	32.8	Quartz Orthoclase Biotite Hornblende	Straw-yellow Almost pure white Light greenish gray Light green	Deep violet Light red Grayish red Carmine-red	I I IV II
GRANITE .....	23.6	Quartz Orthoclase Plagioclase Microcline Biotite Muscovite Hornblende Augite Hyperssthene Tourmaline Zircon Apatite	Gray Bluish gray Bluish gray Bluish gray Bluish gray Greenish yellow Bright orange-red Purple Reddish orange Pale straw-yellow Brownish yellow Greenish yellow Lavender-gray	Brownish yellow Brownish white Brownish white Brownish white Grayish blue Greenish blue Light green Bluish green Dark reddish brown Gray-blue Grayish blue Yellowish white	I I I I III II II I I I III I
PEGMATITE .....	21.4	Quartz Feldspar Biotite Muscovite Tourmaline Beryl	Bluish gray Grayish blue Greenish blue Pure yellow Brownish yellow Grayish blue	Brownish yellow Brownish white Flesh-colored Indigo Gray-blue Brownish white	I I III II I I
TRACHYTYE .....	27.2	Orthoclase Biotite Hornblende Augite	Gray Dull purple Sky-blue Indigo	Brownish yellow Dull sea-green Orange Golden yellow	I III II II
PHONOLITE .....	33.1	Orthoclase Augite Elaeolite	Yellowish white Greenish blue Grayish blue	Carmine-red Brownish orange Brownish white	I II I
TINGUAITE .....	18.4	Orthoclase Biotite Hornblende Augite Leucite Elaeolite	Grayish blue Dark violet-red Bright yellow Brownish yellow Black Lavender-gray	Brownish white Green Blue Gray-blue Bright white Yellowish white	I II I I I I
SYENITE .....	24.7	Orthoclase Augite Hornblende Biotite	Bluish gray Red Indigo Flesh-colored	Brownish yellow Pale green Golden yellow Sea-green	I I II III
NEPHELINE-SYENITE ..	19.6	Orthoclase Nepheline Hornblende Augite Biotite	Grayish blue Lavender-gray Reddish orange Brownish yellow Indigo	Brownish white Yellowish white Bluish green Gray-blue Impure yellow	I I I I III
ANDESITE .....	25.3	Plagioclase Augite Biotite Hornblende	Bluish gray Deep red Carmine-red Indigo	Brownish yellow Yellowish green Green Golden yellow	I I III II

Rock Section.	Thickness in Microns.	Crystals.	Interference Color between Crossed Nicols.	Interference Color between Parallel Nicols.	Order of Color.
DACITE.....	33.7	Quartz Plagioclase Biotite Augite Hornblende	Light yellow Yellowish white Whitish gray Green Light green	Indigo Carmine-red Bluish gray Light carmine-red Purplish red	I I IV II II
DIORITE.....	29.1	Plagioclase Augite Biotite Hornblende	Greenish white Sky-blue Bluish green Greenish blue	Brown Orange Lilac Brownish orange	I II IV II
BASALT.....	38.9	Plagioclase Augite	Light yellow Greenish yellow	Indigo Violet	I II
OLIVINE-BASALT.....	41.7	Plagioclase Augite Olivine	Bright yellow Pure yellow Flesh colored	Blue Indigo Sea-green	I II III
HORNBLENDE-BASALT.	46.3	Plagioclase Augite Hornblende	Bright yellow Bright orange-red Dark violet red	Blue Greenish blue Green	I II II
NEPHELINE-BASALT...	43.6	Nepheline Augite	Gray Orange	Brownish yellow Dark blue	I II
LEUCITE-BASALT.....	45.9	Leucite Augite	Iron-gray Bright orange-red	White Greenish blue	I II
MELILITE-BASALT.....	51.3	Mellilite Augite	Straw-yellow Light bluish violet	Deep violet Yellowish green	I III
DIABASE.....	49.2	Plagioclase Augite	Brownish yellow Dark violet red	Gray-blue Green	I II
GABBRO.....	46.5	Plagioclase Augite	Bright yellow Bright orange-red	Blue Greenish blue	I II
NORITE.....	41.3	Plagioclase Hypersthene	Bright yellow Reddish orange	Blue Bluish green	I I
TROCTOLITE.....	55.4	Plagioclase Olivine	Brownish yellow Whitish gray	Gray-blue Bluish gray	I IV
THERALITE.....	33.4	Augite Biotite Plagioclase Nepheline Apatite	Greenish blue Whitish gray Yellowish white Grayish blue Grayish blue	Brownish orange Bluish gray Carmine-red Brownish white Brownish white	
PERIDOTITE.....	29.8	Olivine Augite Diopside Hornblende Biotite	Dark violet-red Sky-blue Bright orange-red Greenish blue Light green	Green Orange Greenish blue Brownish orange Carmine	II II II II IV

# COLOR-SCALE ACCORDING TO NEWTON.

$\Delta$ IN MICRONS.	INTERFERENCE COLOR BETWEEN CROSSED NICOLS.	INTERFERENCE COLOR BETWEEN PARALLEL NICOLS.
0.000	Black	Bright white
0.040	Iron-gray	White
0.097	Lavender-gray	Yellowish white
0.158	Grayish blue	Brownish white
0.218	Clearer gray	Brownish yellow
0.234	Greenish white	Brown
0.259	Almost pure white	Light red
0.267	Yellowish white	Carmine-red
0.275	Pale straw-yellow	Dark reddish brown
0.281	Straw-yellow	Deep violet
0.306	Light yellow	Indigo
0.332	Bright yellow	Blue
0.430	Brownish yellow	Gray-blue
0.505	Reddish orange	Bluish green
0.536	Red	Pale green
0.551	Deep red	Yellowish green
0.565	Purple	Lighter green
0.575	Violet	Greenish yellow
0.589	Indigo	Golden yellow
0.664	Sky-blue	Orange
0.728	Greenish blue	Brownish orange
0.747	Green	Light carmine-red
0.826	Lighter green	Purplish red
0.843	Yellowish green	Violet-purple
0.866	Greenish yellow	Violet
0.910	Pure yellow	Indigo
0.948	Orange	Dark blue
0.998	Bright orange-red	Greenish blue
1.101	Dark violet-red	Green
1.128	Light bluish violet	Yellowish green
1.151	Indigo	Impure yellow
1.258	Greenish blue	Flesh-colored
1.334	Sea-green	Brownish red
1.376	Brilliant green	Violet
1.426	Greenish yellow	Grayish blue
1.495	Flesh-colored	Sea-green
1.534	Carmine-red	Green
1.621	Dull purple	Dull sea-green
1.652	Violet-gray	Yellowish green
1.682	Grayish blue	Greenish yellow
1.711	Dull sea-green	Yellowish gray
1.744	Bluish green	Lilac
1.811	Light green	Carmine
1.927	Light greenish gray	Grayish red
2.007	Whitish gray	Bluish gray

FIRST ORDER

SECOND ORDER

THIRD ORDER

FOURTH ORDER

## DOUBLY-REFRACTING POWER OF VARIOUS MINERALS.

CRYSTAL.	$n_1 - n_2$	CRYSTAL.	$n_1 - n_2$
Proustite . . . . .	0.2953	Phenacite . . . . .	0.0157
Sulphur . . . . .	0.2900	Anorthite . . . . .	0.0130
Rutile . . . . .	0.2871	Natrolite . . . . .	0.0119
Dolomite . . . . .	0.1791	Hyphersthene . . . .	0.0115
Calcite . . . . .	0.1722	Andalusite . . . . .	0.0110
Brookite . . . . .	0.1580	Gypsum . . . . .	0.0100
Aragonite . . . . .	0.1558	Bronzite . . . . .	0.0100
Cassiterite . . . . .	0.0968	Corundum . . . . .	0.0092
Zircon . . . . .	0.0618	Quartz . . . . .	0.0091
Biotite . . . . .	0.0600	Topaz . . . . .	0.0090
Epidote . . . . .	0.0545	Cordierite . . . . .	0.0090
Anhydrite . . . . .	0.0430	Axinite . . . . .	0.0090
Muscovite . . . . .	0.0420	Orthoclase . . . . .	0.0080
Olivine . . . . .	0.0360	Albite . . . . .	0.0080
Scapolite . . . . .	0.0360	Beryl . . . . .	0.0073
Diopside . . . . .	0.0335	Plagioclase . . . . .	0.0071
Cancrinite . . . . .	0.0289	Zoisite . . . . .	0.0060
Tremolite . . . . .	0.0265	Melilite . . . . .	0.0058
Hornblende . . . . .	0.0240	Nephelite . . . . .	0.0049
Augite . . . . .	0.0220	Elaeolite . . . . .	0.0047
Sillimanite . . . . .	0.0220	Apatite . . . . .	0.0042
Glaucomphane . . . . .	0.0216	Vesuvianite . . . . .	0.0020
Brucite . . . . .	0.0210	Apophyllite . . . . .	0.0019
Tourmaline . . . . .	0.0189	Leucite . . . . .	0.0010
		Chlorite . . . . .	0.0010

---

---

**Plate 2.**

---

---

PLATE 2. (Crossed Nicols.)

Chalcedony,

Section cut at right angles to the length of a stalactite;  $\times 100$  diameters. Central sections through the concretions give a spherulitic interference cross. Radial fibers appear to be uniaxial, since extinction is complete in fibers parallel to planes of vibration of polarizer and analyzer, i. e., when  $\Psi = 0$ ; the illumination is at a maximum when  $\Psi = 45^\circ$ . Extinction is also complete where fibers are cut at right angles, for the light travels in the direction of their optic axes. The crystal system of this mineral has not yet been definitely determined; however, it appears to be optically uniaxial.



DERR



## DERIVATION OF FORMULAE.

Let there be a crystalline plate having OX and OY for axes of elasticity, OP the direction of a polarizer on this section (fig. 5); a ray of light polarized by OP vibrating in the principal section of the polarizer is divided, in the plate, into two rays, which may be regarded as superposed and

which have at emergence a difference of phase :

$$\varphi = \frac{h}{\lambda \cos r} (n_e - n_o).$$

Let us call x and y the vibrations of these rays which give as a resultant an elliptical movement, and let us receive the corresponding doubly-refracted ray on an analyzer having OA for the direction on the plane of the plate; the angle between the polarizer and the axis OX of the plate is POX. Put  $\angle POX = \beta$ . The angle between the analyzer and the same axis is AOX, and we put  $\angle AOX = \alpha$ . The components belonging to OX and OY are :

$$x = a \cos 2 \pi \frac{t}{T} \cos \beta,$$

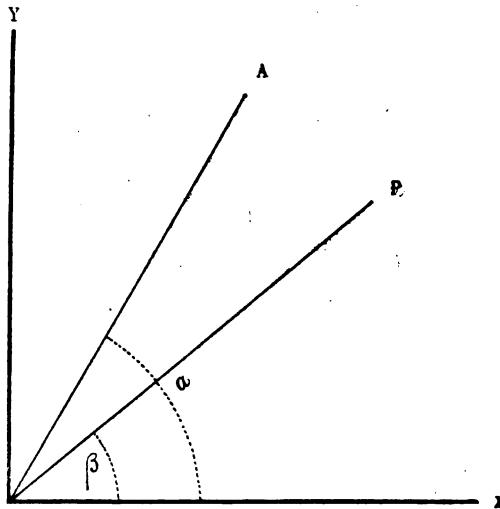
$$y = a \cos 2 \pi \frac{t}{T} \sin \beta.$$

They give rise to two rays which emerge from the plate with a difference of phase  $\varphi$ ; the components of these rays are :

$$x = a \cos 2 \pi \frac{t}{T} \cos \beta, y = a \cos 2 \pi \left[ \frac{t}{T} - \varphi \right] \sin \beta.$$

x gives to OA a component  $\xi = x \cos \alpha$  parallel to the principal section and a perpendicular; y gives to OA a component  $\eta = y \sin \alpha$  parallel to its principal section and a perpendicular to that section. We need not concern ourselves with the perpendiculars, which are eliminated; the parallels give a resultant, which is their algebraic sum :

$$\xi + \eta = x \cos \alpha + y \sin \alpha.$$



Replacing  $x$  and  $y$  by their values, we have

$$\begin{aligned}\xi + \eta &= a \cos \alpha \cos \beta \cos 2 \pi \frac{t}{T} + a \sin \alpha \sin \beta \cos 2 \pi \left( \frac{t}{T} - \varphi \right) \\ &= a \cos 2 \pi \frac{t}{T} (\cos \alpha \cos \beta + \sin \alpha \sin \beta \cos 2 \pi \varphi) \\ &\quad + a \sin 2 \pi \frac{t}{T} \sin \alpha \sin \beta \sin 2 \pi \varphi.\end{aligned}$$

According to Fresnel, let us put

$$\begin{aligned}a (\cos \alpha \cos \beta + \sin \alpha \sin \beta \cos 2 \pi \varphi) &= E \cos \Delta, \\ a \sin \alpha \sin \beta \sin 2 \pi \varphi &= E \sin \Delta.\end{aligned}$$

$$\eta + \xi = \cos 2 \pi \frac{t}{T} \cos \Delta + \sin 2 \pi \frac{t}{T} \sin \Delta = \cos \left[ 2 \pi \frac{t}{T} - \Delta \right].$$

This relation shows that the amplitude of vibration is  $E$ , if we compare it to  $x = a \cos 2 \pi \frac{t}{T}$ .

We have :

$$E^2 (\cos^2 \Delta + \sin^2 \Delta) = E^2 = \left\{ \begin{aligned} &a^2 (\cos \alpha \cos \beta + \sin \alpha \sin \beta \cos 2 \pi \varphi)^2 \\ &+ a^2 \sin^2 \alpha \sin^2 \beta \sin^2 2 \pi \varphi.\end{aligned} \right.$$

$$\begin{aligned}\frac{E^2}{a^2} &= \cos^2 \alpha \cos^2 \beta + \sin^2 \alpha \sin^2 \beta \cos^2 2 \pi \varphi \\ &\quad + 2 \cos \alpha \cos \beta \sin \alpha \sin \beta \cos 2 \pi \varphi + \sin^2 \alpha \sin^2 \beta \sin^2 2 \pi \varphi \\ &= \cos^2 \alpha \cos^2 \beta + \sin^2 \alpha \sin^2 \beta + 2 \sin \alpha \sin \beta \cos \alpha \cos \beta \\ &\quad - 4 \sin \alpha \sin \beta \cos \alpha \cos \beta \sin^2 \pi \varphi.\end{aligned}$$

Replacing  $\cos 2 \pi \varphi$  by  $1 - 2 \sin^2 \pi \varphi$ ,

$$\begin{aligned}\frac{E^2}{a^2} &= (\cos \alpha \cos \beta + \sin \alpha \sin \beta)^2 - \sin 2 \alpha \sin 2 \beta \sin^2 \pi \varphi \\ &= \cos^2 (\alpha - \beta) - \sin 2 \alpha \sin 2 \beta \sin^2 \pi \varphi . . . . . \quad (1)\end{aligned}$$

### DISCUSSION OF THE FORMULA.

We have that  $\varphi = \frac{h}{\lambda \cos r} (n_e - n_o)$ .

$h$  represents the thickness of the section ;  $r$  the angle between the rays, which we regard as superposed, and the normal to the faces of the section ;  $n_e$ ,  $n_o$  the indices of refraction of the two rays.

When white light, composed of an infinite number of lights of very different wave-lengths, is employed, the second term corresponds to the sum of the intensities of these lights, and the general formula may then be written :

$$\frac{E^2}{a^2} = \cos^2 (\alpha - \beta) - \sin 2 \alpha \sin 2 \beta \sum \sin^2 \pi \frac{h}{\lambda \cos r} (n_e - n_o) . . . \quad (2)$$

## THE DIFFERENT CASES.

I.  $\alpha - \beta = \frac{\pi}{2}$ , case where the principal sections of the polarizer and analyzer are at right angles to each other.

Since  $\beta = \alpha - \frac{\pi}{2}$ , formula (2) becomes :

$$\frac{E^2}{a^2} = \cos^2 \frac{\pi}{2} - \sin 2 \alpha \sin (2 \alpha - \pi) \sum \sin^2 \pi \varphi, \text{ or}$$

$$\frac{E^2}{a^2} = \sin^2 2 \alpha \sum \sin^2 \pi \varphi = \sin^2 2 \alpha \sum \sin^2 \left[ \frac{\pi \Delta}{\lambda \cos r} \right] \dots \dots \quad (3)$$

### PARTICULAR CASES.

1°. Let  $\beta = 0$ , then  $\alpha = \frac{\pi}{2}$ . The axes OX, OY of the plate are parallel to the principal sections of both analyzer and polarizer. All the colors are extinguished in the crystal at the same time, since they have the same axes of elasticity.

$$\frac{E^2}{a^2} = 0.$$

2°.  $\beta = 45^\circ$ ,  $\alpha = \frac{\pi}{2} + 45^\circ$ .

$$\frac{E^2}{a^2} = \sum \sin^2 \pi \varphi. \dots \dots \dots \dots \dots \dots \quad (4)$$

II.  $\alpha = \beta$ , the polarizer and analyzer are parallel.

$$\frac{E^2}{a^2} = 1 - \sin^2 2 \alpha \sum \sin^2 \pi \varphi. \dots \dots \dots \dots \dots \dots \quad (5)$$

The values (3) and (5) are complementary.

2°.  $\alpha = 45^\circ$ ,  $\beta = 45^\circ$ .

$$\frac{E^2}{a^2} = 1 - \sum \sin^2 \pi \varphi. \dots \dots \dots \dots \dots \dots \quad (6)$$

All the colors have their maximum of intensity.

---

If we employ a rhombohedron of calcite instead of a nicol, we have two images. The intensity  $\frac{O^2}{a^2}$  of the ordinary image is obtained by means of formula (2) by replacing  $\alpha$  by  $\alpha + \frac{\pi}{2}$ ; the formula becomes :

$$\begin{aligned} \frac{O^2}{a^2} &= \sin^2 \alpha \cos^2 \beta + \sin^2 \beta \cos^2 \alpha \\ &\quad - 2 \sin \alpha \sin \beta \cos \alpha \cos \beta \cos 2 \pi \varphi \\ &= \sin^2(\alpha - \beta) + \sin 2 \alpha \sin 2 \beta \sum \sin^2 \pi \varphi. \dots \quad (7) \end{aligned}$$

These phenomena are complementary to those which give  $\frac{E^2}{a^2}$ .

## BIBLIOGRAPHY.

“Theory of Light,” . . . . . PRESTON.  
“Lehrbuch der Optik,” . . . . . DRUDE.  
“Light,” . . . . . TAIT.  
“Physical Optics,” . . . . . GLAZEBROOK.  
“General Physics,” . . . . . HASTINGS and BEACH.  
“Mikroskopische Physiographie der pet-  
rographisch wichtigen Mineralien, . ROSENBUSCH.







UNIVERSITY OF CALIFORNIA LIBRARY

THIS BOOK IS DUE ON THE LAST DATE  
STAMPED BELOW

SEP 18 1915

3/13/16

FEB 10 1917

FEB 6 1916

30m-1-'15

YC 40243



